A multicomponent model for the optical to γ -ray emission from the Crab Pulsar

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Abstract. We present a multicomponent model to explain the features of the pulsed emission and spectrum of the Crab Pulsar, on the basis of X and γ -ray observations performed with BeppoSAX, INTEGRAL and CGRO. This model explains the evolution of the pulse shape and of the phase-resolved spectra, ranging from the optical/UV to the GeV energy band, on the assumption that the observed emission is due to several components. The first component, C_O , is assumed to have the pulsed double-peaked profile observed at the optical frequencies, while the second component, C_X , is dominant in the interpeak and second peak phase regions. The spectra of these components are modelled with log-parabolic laws. Moreover, to explain the properties of the pulsed emission in the MeV-GeV band, we introduce two more components, $C_{O\gamma}$ and $C_{X\gamma}$, with phase distributions similar to those of C_O and C_X and log-parabolic spectra with the same curvature but different peak energies. This multicomponent model is able to reproduce both the broadband phase-resolved spectral behaviour and the changes of the pulse shape with energy. We also propose some possible physical interpretations in which C_O and C_X are emitted by secondary pairs via a synchrotron mechanism while $C_{O\gamma}$ and $C_{X\gamma}$ can originate either from Compton scattered or primary curvature photons.

1. Introduction

The study of the phase distributions of pulsars' signals in the various bands of the electromagnetic spectrum is important to obtain information on the geometry and location of the emission regions in the magnetosphere. At γ -ray energies, in particular, the three brightest sources (Vela, Crab and Geminga) show remarkably similar patterns with two main peaks at a phase separation ranging from 0.40 to 0.48.

The Crab Pulsar (PSR B0531+21) is characterized by a rather stable phase distribution throughout the whole electromagnetic spectrum with a double peak structure. It is well known that the pulse shape of the Crab changes with energy in the X and soft gamma-ray ranges where the emission of the second peak (P2) becomes higher than the first one (P1), and where it is present a significant emission from the region between the two peaks (bridge or interpeak, IP). This behaviour continues up to about 10 MeV, where the pulse almost sharply returns to a shape similar to the optical light curve. A satisfactory explanation for these changes has not been found so far.

On the basis of high quality BeppoSAX data, covering a wide energy range (from 0.1 to 300 keV), we already proposed a two component model (Massaro *et al.*, 2000) to explain the behaviour of the light curve. Here we extend this model, reanalysing the whole set of BeppoSAX Crab observations with new ISGRI-INTEGRAL data at higher energies (Mineo et al., 2006). We found that the energy spectra of these components are not described by a simple power law, but show a spectral steepening towards high energies. We model these components with log-parabolic spectral distributions. Moreover, to explain the behaviour in the MeV/GeV band as observed by COMPTEL and EGRET onboard Compton-CGRO, two more components are introduced, both with a similar shape and spectrum of the X-ray counterparts. A complete description of the data analysis and of the model can be found in Massaro et al. (2006).

2. The two-component model: optical to hard X-rays

As presented in Massaro *et al.* (2000), Crab X-ray light curve is well reproduced by two phase-components. The first component, called C_O , is assumed to have the same pulsed profile observed at optical frequencies, while the second component, C_X , is described by an analytical model whose shape is determined by comparing $C_O + C_X$ with the observed pulse profiles, and that dominates at the interpeak (IP) and second peak (P2) phase regions (Fig. 1).

Using the high-statistics observations of BeppoSAX we performed a phase-resolved spectral analysis and found that the photon indices of P1, P2 and IP are changing with energy, and linearly increasing with Log E (Fig. 2).



Fig. 1. The two components C_O and C_X of the model at the energies of 8.85 keV (left) and 75.2 keV (right). In the upper panels: the model compared with BeppoSAX data. In the lower panels: C_O and C_X (adapted from Massaro *et al.* 2000).



Fig. 2. Photon indices of P1, IP and P2 as measured by the four NFI of BeppoSAX and by INTEGRAL-ISGRI.

We found that the spectra of C_O and C_X are well fitted by a log-parabolic spectral law,

$$F(E) = KE^{-(a+b\operatorname{Log} E)} \tag{1}$$

where K is the flux at 1 keV and E is the energy in keV. The parameter b describes the "curvature" of the logparabola. The energy-dependent spectral index can be obtained from the previous equation: $\alpha(E) = a + 2b \operatorname{Log} E$. According to this spectral law, the spectral energy distribution (SED) has a maximum at the energy $E_p = 10^{(2-a)/2b}$. The curvature parameter b is equal to 0.16 for both C_O and C_X , while the peak energies are respectively 12 keV and 178 keV.



Fig. 3. Broadband spectra of the total averaged pulse with the four components of the model. Upward pointing triangles: LECS; brown circles: MECS; diamonds: HPGSPC; downward pointing triangles: PDS; leftward pointing triangles: FIGARO II; squares: COMPTEL; crosses: EGRET. Dashed line: C_O ; dash-dotted line: C_X ; dot-dot-dashed line: $C_{O\gamma}$; dash-dash-dotted line: $C_{X\gamma}$.

3. Extension of the model to the MeV/GeV band: the need for two more components

CGRO COMPTEL and EGRET observations (Kuiper et al., 2001; Thompson, 2004) provided above ~ 10 MeV light curves of a good statistical quality which show that the pulse shape is similar to that of C_O , although some minor differences are present. At energies higher than $\sim 500 \text{ MeV}$ the emission from IP and P2 increases, and this seems to reproduce the behaviour of the X-ray emission. In order to explain such a finding, we assume that there are two more, high-energy spectral components, $C_{O\gamma}$ and $C_{X\gamma}$, both with a log-parabolic spectral distribution and with the same pulse shape of the lower-energy components C_{O} and C_X . To be consistent with the upper limits to the TeV pulsed emission (e.g. Lessard *et al.*, 2000) we added also an exponential cutoff to both $C_{O\gamma}$ and $C_{X\gamma}$, at the energy $E_c = 15$ GeV. This model therefore has 6 adjustable parameters, i.e. the peak energies, curvatures and normalizations of the $C_{O\gamma}$ and $C_{X\gamma}$ components. Assuming that the curvatures are equal to the C_O and C_X ones (b = 0.16), we are then able to reproduce the broadband energy spectrum of the total (averaged) pulse and of the P1, IP and P2 phase regions (see Figs. 3, 4 and 5) and the ratios of P2/P1 and IP/P1 fluxes (in the same phase intervals of Kuiper et al., 2001; Figs. 6 and 7). We stress that there is no constraint on E_c : in fig. 6 we plot also the P2/P1 ratio for various values of $C_{O\gamma}$ cutoff energy ranging from 9 to $15 \, \mathrm{GeV}.$



Fig. 4. Broadband spectra of P1 with the four components of the model. Upward pointing triangles: LECS; brown circles: MECS; diamonds: HPGSPC; downward pointing triangles: PDS; leftward pointing triangles: IS-GRI; squares: COMPTEL; crosses: EGRET. Dashed line: C_O ; dash-dotted line: C_X ; dot-dot-dashed line: $C_{O\gamma}$; dashdash-dotted line: $C_{X\gamma}$.



Fig. 5. Broadband spectra of P2 with the four components of the model. Upward pointing triangles: LECS; brown circles: MECS; diamonds: HPGSPC; downward pointing triangles: PDS; leftward pointing triangles: IS-GRI; squares: COMPTEL; crosses: EGRET. Dashed line: C_O ; dash-dotted line: C_X ; dot-dot-dashed line: $C_{O\gamma}$; dashdash-dotted line: $C_{X\gamma}$.

4. Physical interpretation

An open question is the physical origin of these components, that phenomenologically explain the observations with a very good approximation, in the framework of the



Fig. 6. P2/P1 ratio as derived from the model. Data points come from various experiments (Kuiper et al 2001). The various extrapolations above 1 GeV correspond to different values of the cut-off energy of the $C_{O\gamma}$ spectrum; from top to bottom: 9, 11, 13 and 15 GeV.



Fig. 7. IP/P1 ratio as derived from the model. Data points come from various experiments (Kuiper et al 2001).

high-energy pulsar emission models, either in the polar cap or outer gap models (e.g. Cheng *et al.*, 2000; Zhang & Cheng, 2002).

Assuming that the lower-energy components C_O and C_X are produced by synchrotron emission of secondary electron-positron pairs created in the pulsar magneto-sphere, the higher-energy components $C_{O\gamma}$ and $C_{X\gamma}$ could be due to:

1. Emission of curvature radiation from primary particles accelerated in the magnetospheric gaps.

2. Emission from inverse Compton scattering of the synchrotron photons by the secondary pairs themselves (Synchrotron-Self-Compton mechanism).

The different shape of the "O" and the "X" components is presumably due to the different location in the magnetosphere of the emission regions.

5. Conclusions

Several models have appeared in the literature based on either polar cap or outer gap geometries. Usually, these models are focused on reproducing either the total spectrum or the phase profile, and generally they are not fully satisfactory in explaining the complex observational picture. Moreover, the possibility that the observed features of the pulsed signal can arise from the superposition of two or more distinct components is not taken into account.

We followed another approach and searched for a possible interpretation of the Crab signal based on the superposition of two or more components that provides a consistent description of the spectral and phase distributions. Clearly, it is only a phenomenological model, but it could furnish some constraints to more detailed, physically-based emission models. In particular it is important to verify whether at energies higher than ~1 GeV the pulse shape tends to be dominated by $C_{X\gamma}$. The GLAST/LAT experiment (Gehrels *et al.*, 1999), with its large collecting area, will give us very useful data in this range that will permit to better estimate the model parameters.

Another interesting perspective is whether this model can be adapted to the other γ -ray pulsars. For Vela and Geminga the main problem is that their pulse profiles change very much in different spectral bands and no clear trend, like the P2/P1 ratio in Crab, has been found to now. In the γ -ray band, however, Kanbach (1999) showed that the peak ratios of all these three pulsars have a rather similar behaviour. This can be an indication that geometrical effects may be more relevant at energies lower than γ -rays and that components like C_O or C_X , if existing in Vela and Geminga, are not detected because they are more beamed than the high energy photons.

Acknowledgements. This work was financially supported by Università di Roma La Sapienza. R.C. also acknowledges the support by the WE-Heraeus foundation during the seminar.

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